

# Prospects for little Higgs models at the LHC

Eduardo Ros

IFIC-Valencia

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**Abstract.** The ATLAS Collaboration at the LHC is presently investigating the possibility to detect particles predicted by Little Higgs models. In this talk, the possibility to detect the heavy gauge boson  $Z_H$  and its subsequent decay into  $Zh$  is reviewed.

## 1 Introduction

Little Higgs models have been recently proposed as a possible solution to the hierarchy problem. They try to explain the smallness of the Higgs boson mass by introducing new particles at the 1 TeV scale. In the so-called ‘littlest Higgs model’ [1], these new particles are scalars ( $\Phi^0$ ,  $\Phi^+$ ,  $\Phi^{++}$ ), gauge bosons ( $W_H$ ,  $Z_H$ ,  $A_H$ ) and a heavy top quark ( $T$ ). The masses and couplings of all these new particles are completely fixed (except for  $A_H$ ) once the scale  $f$  and a set of couplings called  $v'$ ,  $\theta$ ,  $\theta'$  and  $\lambda_1$  are specified.

## 2 Phenomenology at the LHC

Branching ratios and cross-sections at the LHC have been computed in [2]. From these calculations it is possible to extract some conclusions concerning experimental strategies in order to test Little Higgs models with LHC experiments:

- The scalar  $\Phi^{++}$  is produced in  $W^+W^+$  fusion (VBF mechanism). The cross-section is proportional to  $(v'/v)^2$ , where  $v = 244$  GeV is the Fermi scale, but  $v'/v$  is expected to be small. The dominant decay mode would be  $\Phi^{++} \rightarrow W^+W^+$ , and the Standard Model background for this mode is rather large, so this particle is difficult to observe.
- The heavy top quark  $T$  is produced according to  $bq \rightarrow Tq'$  via  $W$  exchange in the t-channel ( $Wb$  fusion mechanism). The cross-section is proportional to  $\lambda_1^2$  and  $\lambda_1$  is expected to be of order 1, but the  $b$ -quark content of the proton is small so the cross-section is also small. Therefore this particle is also difficult to observe.
- The gauge boson  $Z_H$  is produced in  $q\bar{q}$  annihilation, in the same way as a normal  $Z$  boson. The cross-section is proportional to  $(\cot\theta)^2$ , the mixing angle  $\theta$  being the only free parameter of the theory once the mass of  $Z_H$  is fixed. For a mass of 2 TeV and  $\cot\theta = 1$ , the cross-section is 1 pb, so  $Z_H$  is copiously produced at the LHC. The charged gauge boson  $W_H$  is also produced

in  $q\bar{q}'$  collisions, the same as the  $W$ , and the production cross-section is also large. The production cross-section for the gauge boson  $A_H$  is more difficult to calculate since the couplings are not entirely fixed by the model.

In the following we concentrate in the experimental search for  $Z_H$  using the ATLAS experiment at the LHC.

## 3 Experimental search for $Z_H$

Once the heavy gauge boson  $Z_H$  is produced, it decays into quark or lepton pairs. Taking into account the universality of the coupling and neglecting fermion masses,  $\text{BR}(Z_H \rightarrow l^+l^-) = 1/24 = 4.2\%$ , where  $l$  is any charged lepton. At small values of  $\cot\theta$ , however, the decay  $Z_H \rightarrow Zh$ , where  $h$  is the Higgs boson, is dominant.

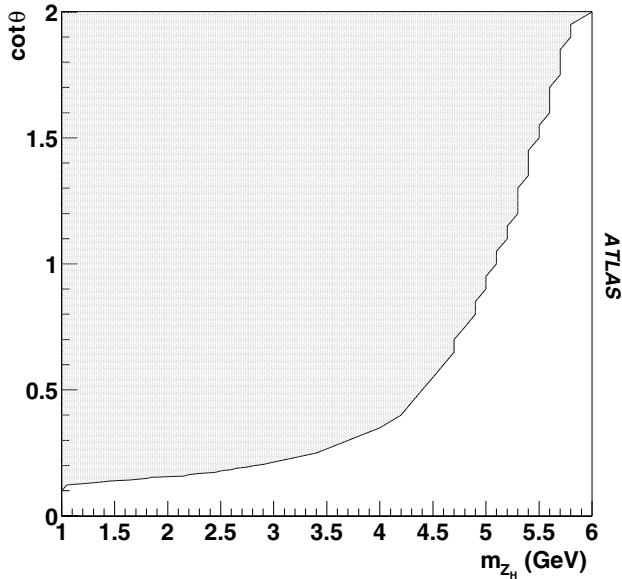
The decay  $Z_H \rightarrow e^+e^-$  provides the best signature to detect  $Z_H$  at the LHC, since the relative invariant mass resolution,  $\sigma(M)/M$ , does not degrade for electrons with increasing mass (contrary to the muon case). The background, mainly Drell-Yan pairs, is much smaller than the signal over a wide range of  $\cot\theta$  values. Fig.1 shows the region in the  $M - \cot\theta$  plane where a discovery of  $Z_H$  is possible at the LHC, using the ATLAS detector. A luminosity of  $3 \cdot 10^5 \text{ pb}^{-1}$  corresponding to 3 years of running at high luminosity has been assumed. If a signal is detected, it would also be possible to measure  $\cot\theta$  via the cross-section and the width of  $Z_H$ . Indeed, as mentioned before, the cross-section is proportional to  $(\cot\theta)^2$ , and the total width of  $Z_H$  is:

$$\Gamma/M = [3.4(\cot\theta)^2 + 0.071(\cot2\theta)^2]\%$$

In this expression the first term accounts for the decay of  $Z_H$  into fermions and the second term for the decay  $Z_H \rightarrow Zh$ .

## 4 Search for the decay $Z_H \rightarrow Zh$

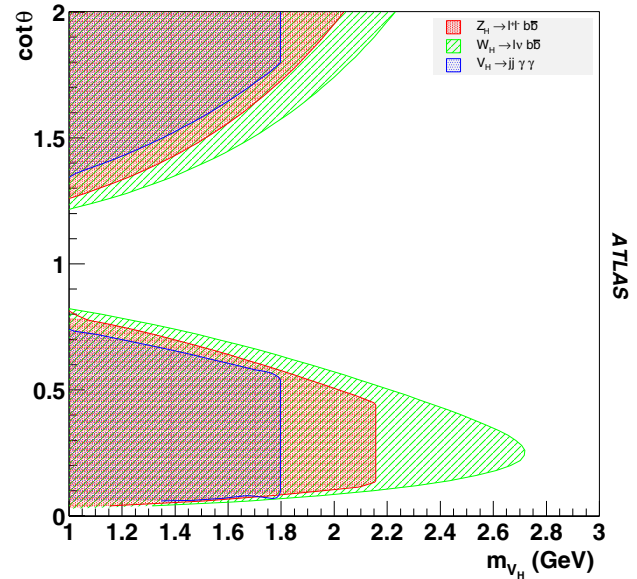
The observation of the decay  $Z_H \rightarrow Zh$  is essential to test Little Higgs models. The amplitude of the decay is not



**Fig. 1.** Region where a discovery of the decay  $Z_H \rightarrow e^+e^-$  is possible using the ATLAS detector at the LHC. An integrated luminosity of  $3 \cdot 10^5 \text{ pb}^{-1}$  has been assumed. The discovery region corresponds to a significance of the signal larger than 5.

proportional to  $(\cot\theta)^2$  as for fermions, but to  $(\cot 2\theta)^2$ . Unfortunately, this  $(\cot 2\theta)^2$  factor in the branching ratio has a tendency to cancel the  $(\cot\theta)^2$  factor in the  $Z_H$  production cross-section. As a result, the  $Zh$  event yield is rather small, except for  $\cot\theta$  values around 0.3. In particular, the event yield completely vanishes for both  $\cot\theta = 0$  and  $\cot\theta = 1$ .

The experimental signature of  $Zh$  events depends on the mass of the Higgs boson. In the following  $M(h)=120$  GeV is assumed, so the dominant decay of the Higgs boson is  $h \rightarrow b\bar{b}$ . In this case the final state consists of a pair of  $b$ -jets and a pair of leptons from the  $Z$  decay. The main background is  $Z$  production in association with jets. As the mass of  $Z_H$  increases, the two  $b$ -jets from the  $h$  decay have a tendency to merge into a double  $b$ -quark jet with very high  $p_T$ . In order to identify these events,  $b$ -tagging at high  $p_T$  is therefore extremely important. The result of full-detector simulation studies show that it would be possible to tag these double  $b$ -quark jets with very high  $p_T$ , by simply requiring a reduced tagging efficiency (40% instead of the usual 50%). Fig.2 shows the region in the  $M - \cot\theta$  plane where the decay  $Z_H \rightarrow Zh$  can be detected at the LHC, using the ATLAS detector and a luminosity of  $3 \cdot 10^5 \text{ pb}^{-1}$ . The figure includes also the decay  $b \rightarrow \gamma\gamma$  and the result obtained using the decay  $W_H \rightarrow Wh$ , assuming that  $W_H$  and  $Z_H$  have exactly the same mass. If the mass of the Higgs boson is larger than 120 GeV, the decays into  $W^+W^-$  and  $ZZ$  have to be considered as well.



**Fig. 2.** Region where a discovery of the decay  $Z_H \rightarrow Zh$  is possible using the ATLAS detector at the LHC. The Higgs boson  $h$  is assumed to decay into either  $b\bar{b}$  or  $\gamma\gamma$ . An integrated luminosity of  $3 \cdot 10^5 \text{ pb}^{-1}$  has been assumed. The discovery region corresponds to a significance of the signal larger than 5. The decay  $W_H \rightarrow Wh$  is considered as well.  $V_H$  is either  $Z_H$  or  $W_H$ .

## 5 Summary and outlook

A short summary of the prospects for detecting at the LHC the heavy gauge boson  $Z_H$ , predicted by the Little Higgs model, has been presented. The possibility to detect the decay  $Z_H \rightarrow Zh$  is discussed as well, assuming that the mass of the Higgs boson is 120 GeV. Work concerning the production and decay of other particles predicted by the model, namely  $A_H$ ,  $W_H$ ,  $T$  and  $\Phi^{++}$ , is in progress. Other Little Higgs models including two Higgs doublets, rather than just one, will be considered as well.

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